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BaBar $|V_{xb}|$ and $|V_{tx}|$ workshop
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$B \rightarrow K^*\gamma$ and $B \rightarrow \rho\gamma$ at NLO

from QCD Factorization

in collaboration with Gerhard Buchalla

- I. Introduction
- II. Basic Formulas
- III. Results for $B \rightarrow K^*\gamma$ and $B \rightarrow \rho\gamma$
- IV. Summary

Introduction

radiative FCNC $b \rightarrow s(d)\gamma$ transition

- high sensitivity to New Physics
- large impact of SD QCD corrections
- experimentally accessible already at present

inclusive mode

- $B(B \rightarrow X_s\gamma)_{exp} = (3.41 \pm 0.36) \cdot 10^{-4}$
ALEPH, BABAR, BELLE, CLEO
- HQE \rightarrow perturbative calculation
Adel, Yao
Chetyrkin, Misiak, Münz
Greub, Hurth, Wyler

exclusive mode

Greub, Simma, Wyler
Asatryan, Asatrian, Wyler
Grinstein, Pirjol

- $B(B^0 \rightarrow K^{*0}\gamma)_{exp} = (4.59 \pm 0.55) \cdot 10^{-5}$
 $B(B^+ \rightarrow K^{*+}\gamma)_{exp} = (3.82 \pm 0.78) \cdot 10^{-5}$
 $B(B^0 \rightarrow \rho^0\gamma)_{exp} < 1.5 \cdot 10^{-6}$ at 90% C.L.
 $B(B^+ \rightarrow \rho^+\gamma)_{exp} < 2.8 \cdot 10^{-6}$
CLEO, BELLE, BABAR

- bound state effects essential

→ nonperturbative hadronic form-factors

→ QCD factorization

Beneke, Buchalla, Neubert, Sachrajda

→ systematic model-independent NLL framework

Beneke, Feldmann, Seidel
SWB, Buchalla

Basic Formulas for $B \rightarrow V\gamma$

Beneke, Feldmann, Seidel
 Ali, Parkhomenko
 SWB, Buchalla

$$\mathcal{H}_{eff} = \frac{G_F}{\sqrt{2}} \sum_{p,i} \lambda_p^{\text{CKM}} C_i(\mu) Q_i^p$$

$$Q_{CC}^p = (\bar{s} p)_{V-A} (\bar{p} b)_{V-A}$$

$$Q_{\text{open}} = (\bar{s} p)_{V-A} \sum_q (\bar{q} q)_{V\mp A}$$

$$Q_7 = \frac{e}{8\pi^2} m_b \bar{s} \sigma^{\mu\nu} (1 + \gamma_5) b F_{\mu\nu}$$

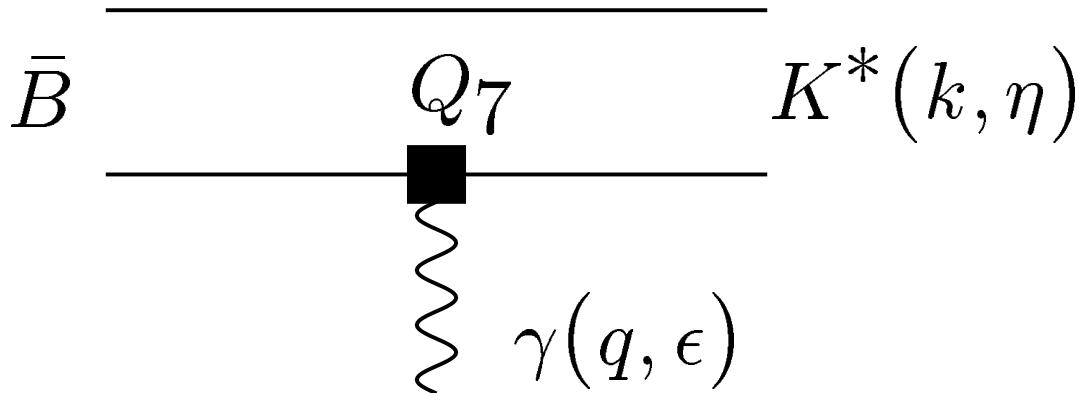
$$Q_8 = \frac{g}{8\pi^2} m_b \bar{s} \sigma^{\mu\nu} (1 + \gamma_5) T^a b G_{\mu\nu}^a$$

matrix elements $\langle V\gamma(\epsilon) | Q_i | \bar{B} \rangle$ in heavy quark limit
 $m_b \gg \Lambda_{QCD} \rightarrow$ factorization formula in 1dg. power

$$\langle Q_i \rangle = \left[F^{B \rightarrow V} T_i^I + \int_0^1 d\xi dv T_i^{II}(\xi, v) \Phi_B(\xi) \Phi_V(v) \right] \cdot \epsilon$$

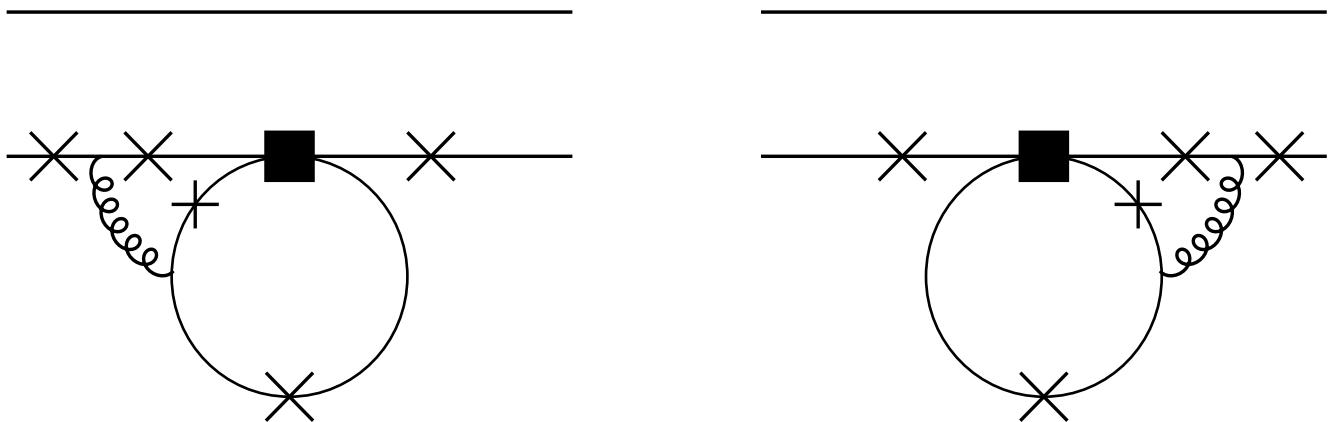
universal nonperturbative form factor and LCDAs
 factorized from perturbative hard-scattering kernels

LO contribution from Q_7



$$\langle Q_7 \rangle = \frac{-e}{2\pi^2} m_b F_V [\varepsilon(\epsilon, \eta, k, q) + i(\epsilon \cdot \eta k \cdot q - \epsilon \cdot k \eta \cdot q)]$$

NLO type I

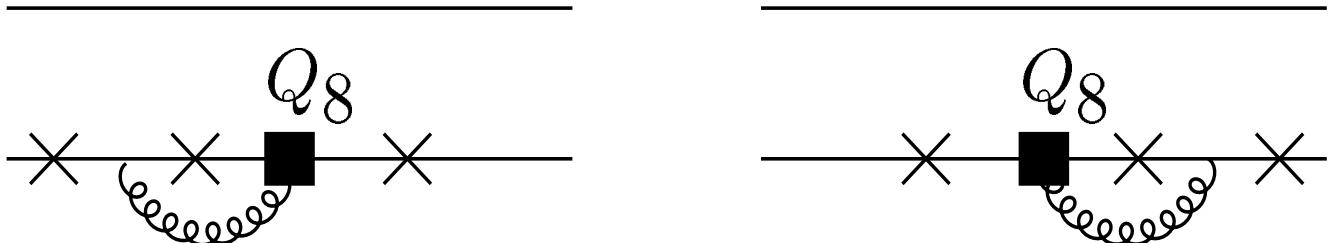


$$\langle Q_i \rangle^I = \langle Q_7 \rangle \frac{\alpha_s C_F}{4\pi} G_i$$

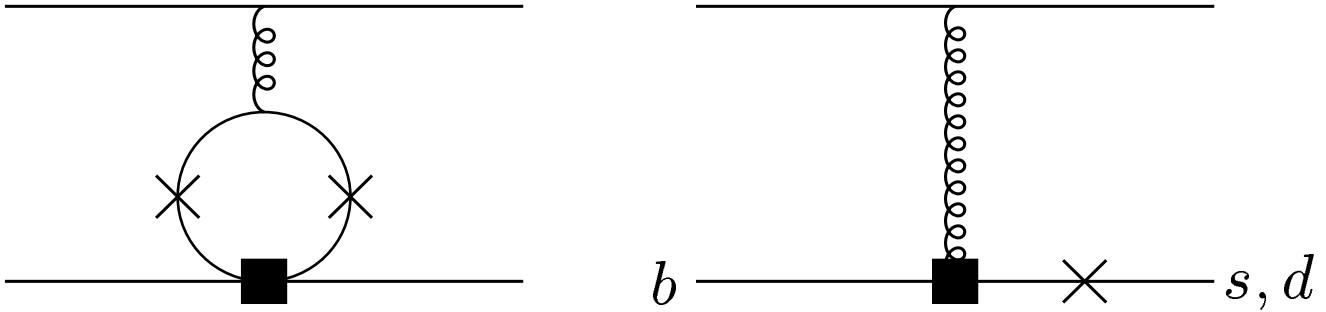
$$G_i(s_c) = l_i \ln \frac{\mu}{m_b} + g_i(\underbrace{s_c}_{=m_c^2/m_b^2})$$

Greub, Hurth, Wyler
Buras, Czarnecki, Misiak, Urban

Dominated by hard scales $\sim m_b \rightarrow$ IR finite

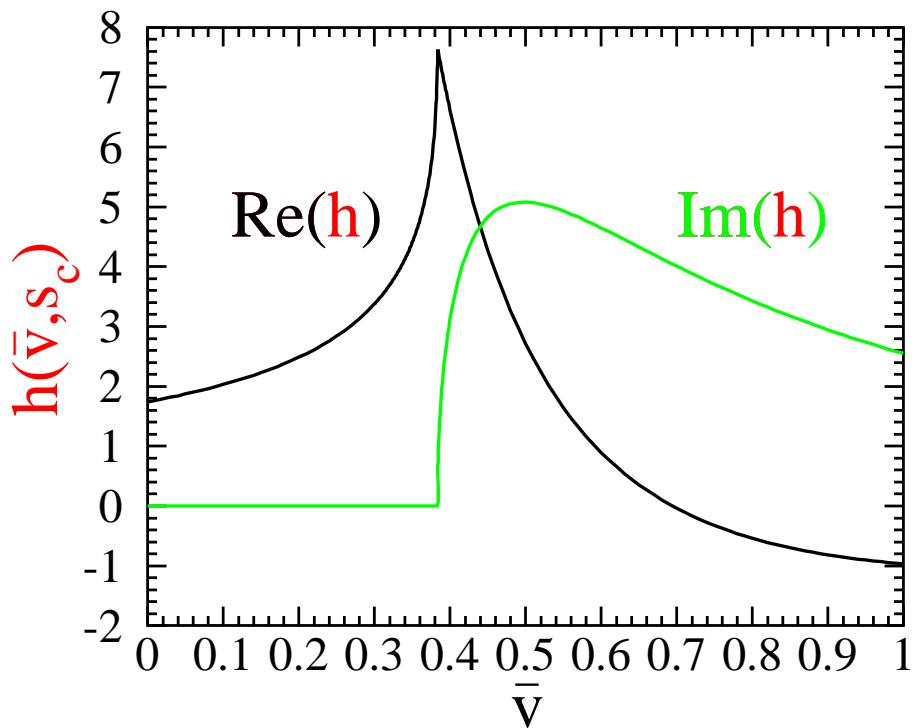


NLO type II



$$\langle Q_i \rangle^{II} = \langle Q_7 \rangle \frac{\alpha_s(\mu_h) C_F}{4\pi} H_i$$

$$H_1^V(s) = -\frac{2\pi^2}{3N} \frac{f_B f_V^\perp}{F_V m_B^2} \underbrace{\int_0^1 d\xi \frac{\Phi_{B1}(\xi)}{\xi}}_{=:m_B/\lambda_B} \int_0^1 dv h(\bar{v}, s) \Phi_\perp^V(v)$$



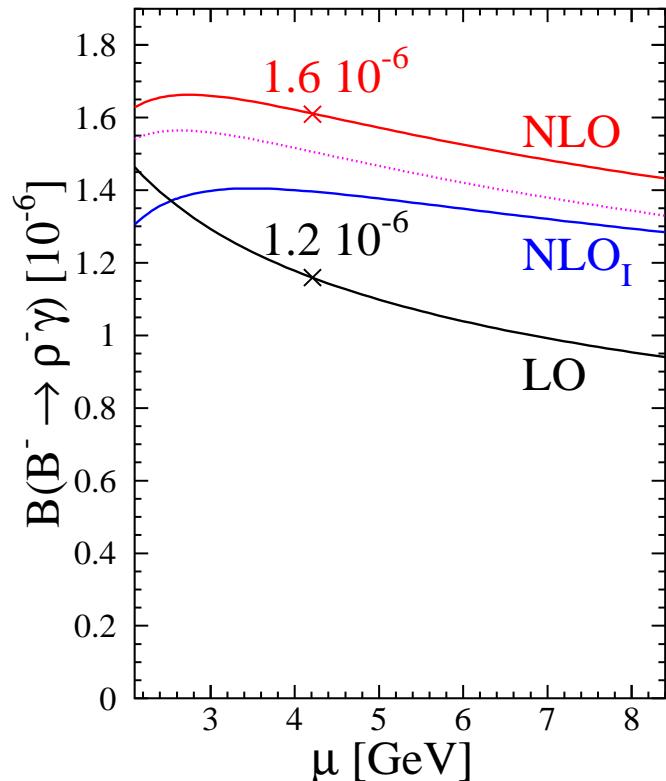
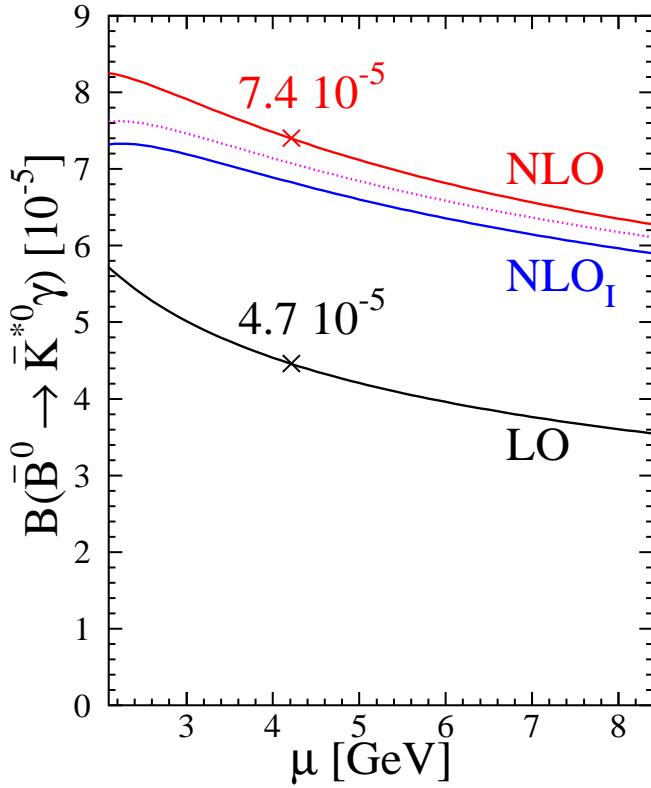
$$H_8^V = \frac{4\pi^2}{3N} \frac{f_B f_V^\perp}{F_V m_B^2} \int_0^1 d\xi \frac{\Phi_{B1}(\xi)}{\xi} \underbrace{\int_0^1 dv \frac{\Phi_\perp^V(v)}{v}}_{=3(1-\alpha_1^\perp + \alpha_2^\perp)}$$

Results

$$A(\bar{B} \rightarrow V\gamma) = \frac{G_F}{\sqrt{2}} \left[\sum_{p=u,c} \lambda_p^{\text{CKM}} a_7^p \right] \langle V\gamma | Q_7 | \bar{B} \rangle$$

$$\begin{aligned} a_7^p(V\gamma) &= C_7 + \frac{\alpha_s(\mu)C_F}{4\pi} \sum_{i=1}^8 C_i(\mu) G_i(s) \\ &\quad + \frac{\alpha_s(\mu_h)C_F}{4\pi} \sum_{j=1}^8 C_j(\mu_h) H_j^V \end{aligned}$$

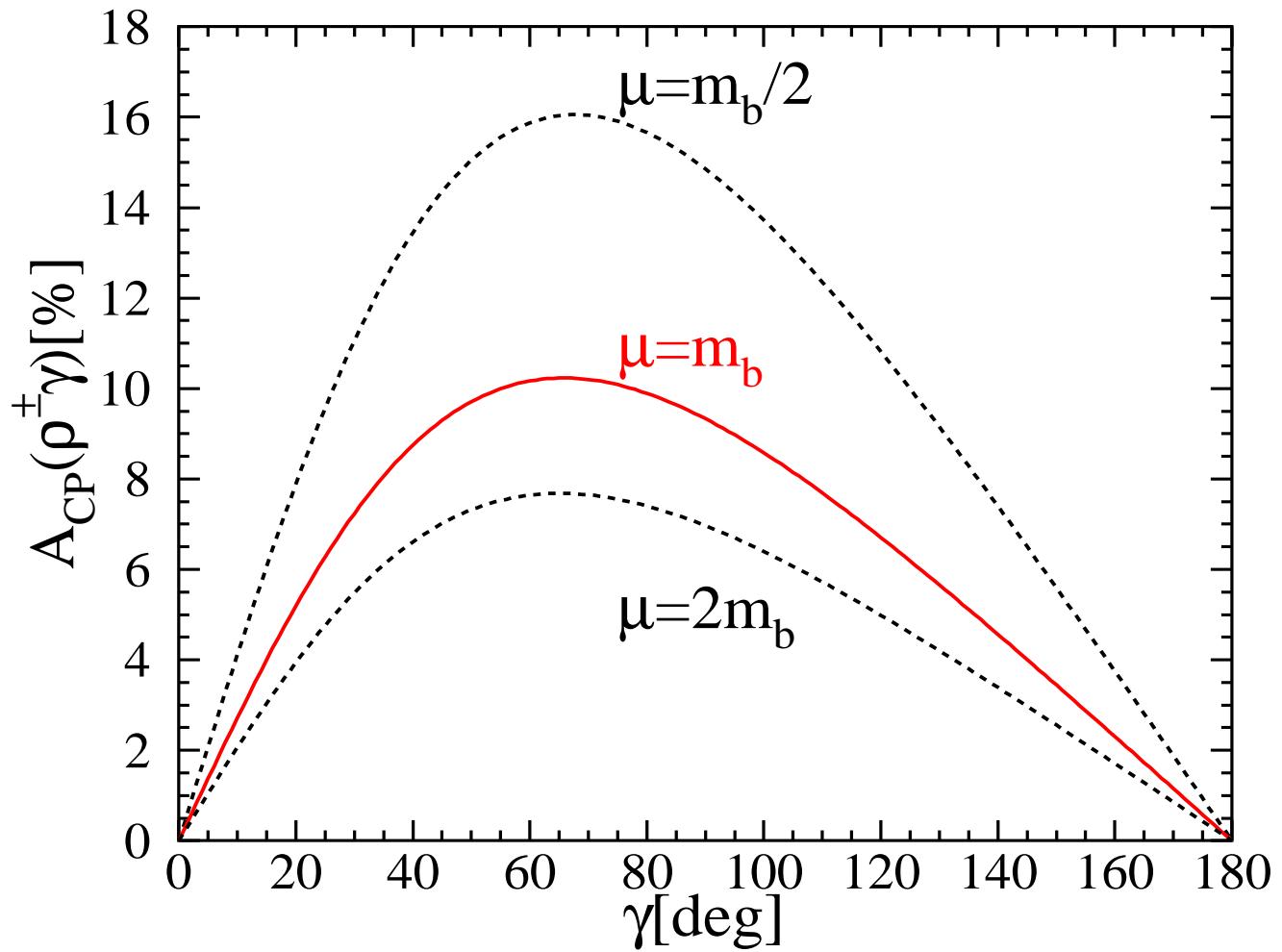
$$\begin{aligned} a_7^c(K^*\gamma) &= -0.322 + 0.011 \\ &\quad -0.079 - 0.013i \\ &\quad -0.017 - 0.012i \\ &= -0.407 - 0.025i \end{aligned}$$



$B \rightarrow \rho\gamma$ Phenomenology

CP asymmetry

$$\begin{aligned} \mathcal{A}_{CP}(\rho\gamma) &= \frac{\Gamma(B \rightarrow \rho\gamma) - \Gamma(\bar{B} \rightarrow \rho\gamma)}{\Gamma(B \rightarrow \rho\gamma) + \Gamma(\bar{B} \rightarrow \rho\gamma)} \\ &\approx \frac{2 \operatorname{Im} \lambda_u^* \lambda_c}{|\lambda_t|^2} \frac{\operatorname{Im} a_7^{u*} a_7^c}{|C_7|^2} \end{aligned}$$



Maximum at $\gamma \approx 65$ deg.

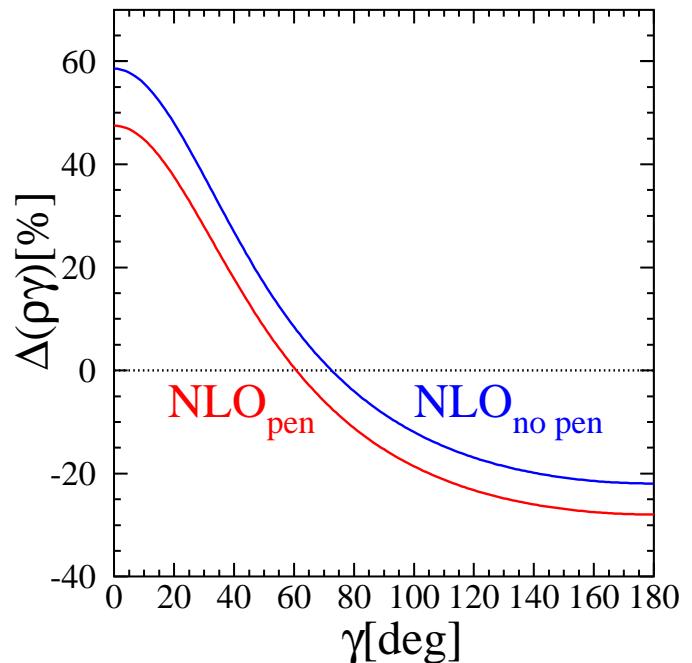
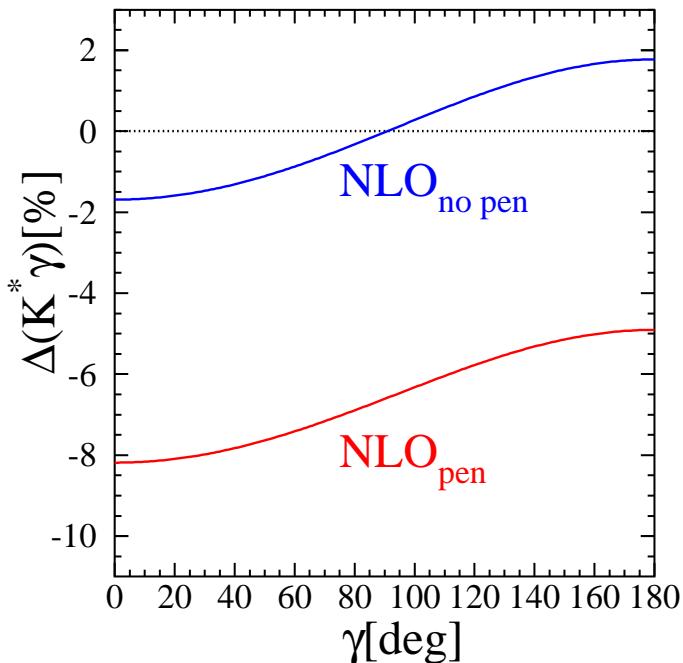
$$\mathcal{A}_{CP}(K^{*0}\gamma) = -0.3\%$$

Isospin breaking

$$\begin{aligned}\Delta_{+0} &= \frac{\Gamma(B^+ \rightarrow \rho^+\gamma)}{2\Gamma(B^0 \rightarrow \rho^0\gamma)} - 1 \\ \Delta_{-0} &= \frac{\Gamma(B^- \rightarrow \rho^-\gamma)}{2\Gamma(\bar{B}^0 \rightarrow \rho^0\gamma)} - 1 \\ \Delta(V\gamma) &= \frac{\Delta_{+0} + \Delta_{-0}}{2}\end{aligned}$$

Ali, Handoko, London
SWB, Buchalla

large effect from Q_6 on $\Delta(K^*\gamma)$ Kagan, Neubert



$$\Delta(K^*\gamma) = (-7.5^{+4.1}_{-5.9})\% \quad \Delta(\rho\gamma) = (2.0^{+27.0}_{-15.7})\%$$

$$\Delta(K^*\gamma)_{\text{exp}} = (-19.2 \pm 11.8)\%$$

Summary

- systematic and model-independent framework for rare radiative decays $B \rightarrow V\gamma$ and $B \rightarrow \gamma\gamma$
- power counting in the heavy quark limit
- quark-loop contributions calculable in QCD factorization
- non-factorizable LD corrections power-suppressed
- strong phases from spectator interaction important for CP asymmetry
- weak annihilation in $B \rightarrow V\gamma$ power suppressed but numerically enhanced and calculable
- NLL: $B(\bar{B}^0 \rightarrow \bar{K}^{*0}\gamma) = (7.4_{-2.4}^{+2.6}) \cdot 10^{-5}$
 $B(B^- \rightarrow \rho^-\gamma) = (1.6_{-0.5}^{+0.7}) \cdot 10^{-6}$
 $A_{CP}(\rho^\pm\gamma) \approx 10\%$
LL: $B(\bar{B}_s \rightarrow \gamma\gamma) = (1.2_{-0.7}^{+2.5}) \cdot 10^{-6}$
 $B(\bar{B}_d \rightarrow \gamma\gamma) = (3.1_{-2.1}^{+6.7}) \cdot 10^{-8}$
- dominant uncertainty: non-perturbative input parameters (form factors, decay constants, λ_B)